

# Statistical Modeling by Wavelets

## 3 Wavelets

### 3.2 Discretization Of The Continuous Wavelet Transform

### 3.3 Multiresolution Analysis

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## Discretization Of The Continuous Wavelet Transform

The critical sampling defined by  $a = 2^{-j}$ ,  $b = k 2^{-j}$ ,  $j, k \in \mathbb{Z}$

For more general sampling, given by

$$a = a_0^{-j}, b = k b_0 a_0^{-j}, j, k \in \mathbb{Z}, a_0 > 1, b_0 > 0$$

numerically stable reconstructions are possible if the system

$\{\psi_{jk}(x), j, k \in \mathbb{Z}\}$  constitutes a frame. Here

$$\psi_{j,k}(x) = a_0^{j/2} \psi\left(\frac{x - k b_0 a_0^{-j}}{a_0^{-j}}\right) = a_0^{j/2} \psi(a_0^j x - k b_0)$$

$$\Phi(w) \text{ in } L_2(\mathbb{R}) \xrightarrow{\text{Fourier Transformation}} \phi(x) \xrightarrow{\text{check orthogonal}} \int \phi(x)\phi(x-k)dx = \delta_k$$

↓ *check orthonormal*

$$\sum |\Phi(w + 2\pi l)|^2 = 1$$

$$\phi(x) \text{ yes no } \mapsto \phi(x) = F^{-1} \left[ \frac{\Phi(w)}{\sqrt{\sum |\Phi(w + 2\pi l)|^2}} \right]$$

↓

$$m_1(w) = -e^{iw} \overline{m_0(w + \pi)} \longleftarrow m_0(w) = \frac{\Phi(w)}{\Phi(w/2)} \longleftarrow \Phi(w)$$

↓

$$\Psi(w) = \Psi\left(\frac{w}{2}\right) m_1(w) \xrightarrow{\text{Fourier transformation}} \psi(x) \text{ in } W_0$$

## Multiresolution Analysis

A multiresolution analysis(MRA) is a sequence of closed subspaces  $V_n, n \in \mathbb{Z}$  in  $L_2(\mathbb{R})$  such that they lie in a containment hierarchy

$$\cdots V_{-2} \subset V_{-1} \subset V_0 \subset V_1 \subset V_2 \subset \cdots \quad (3.8)$$

$$\bigcap_j V_j = \{0\} \quad , \quad \overline{\bigcup_j V_j} = L_2(\mathbb{R})$$

The hierarchy (3.8) is constructed such that

- (1) V-spaces are self-similar,  $f(2^j x) \in V_j$  iff  $f(x) \in V_0$
- (2) There exists a scaling function  $\phi \in V_0$  whose integer-translates span the space  $V_0$ ,

$$V_0 = \left\{ f \in L_2(\mathbb{R}) \mid f(x) = \sum_k c_k \phi(x-k) \right\}$$

and for which the set  $\{\phi(\bullet - k), k \in \mathbb{Z}\}$  is an orthonormal basis.

## Scaling equation

Assume  $\int \phi(x) dx \neq 0$

Since  $V_0 \subset V_1$ , the function  $\phi(x) \in V_0$  can be represented as a linear combination of functions from  $V_1$

$$\text{ie. } \phi(x) = \sum_{k \in \mathbb{Z}} h_k \sqrt{2} \phi(2x - k) \quad (3.10)$$

for some coefficients  $h_k, k \in \mathbb{Z}$

The (possibly infinite) vector  $\tilde{h} = \{h_n, n \in \mathbb{Z}\}$  will be called a **wavelet filter**.

**Note:**

(1)  $\cdots \subset V_2 \subset V_1 \subset V_0 \subset V_{-1} \subset V_{-2} \subset \cdots$  called Daubechies' convention  
and  $\{\phi_{jk}(x) = 2^{-j/2} \phi(2^{-j}x - k), j \text{ fixed}, k \in \mathbb{Z}\}$  is a basis of  $V_j$

(2)  $\cdots \subset V_{-2} \subset V_{-1} \subset V_0 \subset V_1 \subset V_2 \subset \cdots$  called Mallat's convention  
and  $\{\phi_{jk}(x) = 2^{j/2} \phi(2^j x - k), j \text{ fixed}, k \in \mathbb{Z}\}$  is a basis of  $V_j$

Define  $m_0(\omega) = \frac{1}{\sqrt{2}} \sum_{k \in \mathbb{Z}} h_k e^{-ik\omega} = \frac{1}{\sqrt{2}} H(\omega)$  (3.12)

is called **transfer function** and it describes the behavior of the associated filter  $\tilde{h}$  in the Fourier domain.

The period is  $2\pi$  and that the filter taps  $\{h_n, n \in \mathbb{Z}\}$  are the Fourier coefficients of the function  $H(\omega) = \sqrt{2} m_0(\omega)$

Prove  $\Phi(w) = m_0\left(\frac{w}{2}\right) \Phi\left(\frac{w}{2}\right) = \prod_{n=1}^{\infty} m_0\left(\frac{w}{2^n}\right)$

*Pf :*

$$\begin{aligned}
 \Phi(w) &= \int_{-\infty}^{\infty} \phi(x) e^{-iwx} dx \\
 &= \sum_k \sqrt{2} h_k \int_{-\infty}^{\infty} \phi(2x-k) e^{-iwx} dx \\
 &= \sum_k \frac{h_k}{\sqrt{2}} e^{-ikw/2} \int_{-\infty}^{\infty} \phi(2x-k) e^{-i(2x-k)w/2} d(2x-k) \\
 &= \sum_k \frac{h_k}{\sqrt{2}} e^{-ikw/2} \Phi\left(\frac{w}{2}\right) \\
 &= m_0\left(\frac{w}{2}\right) \Phi\left(\frac{w}{2}\right) = \prod_{n=1}^{\infty} m_0\left(\frac{w}{2^n}\right) \quad (\because \Phi(0) = 1)
 \end{aligned}$$



## Normalization

$$\sum_{k \in \mathbb{Z}} h_k = \sqrt{2} \quad (3.16)$$

*pf* :

$$\begin{aligned} \int \phi(x) dx &= \sqrt{2} \sum_k h_k \int \phi(2x - k) dx \\ &= \sqrt{2} \sum_k h_k \frac{1}{2} \int \phi(2x - k) d(2x - k) \\ &= \frac{\sqrt{2}}{2} \sum_k h_k \int \phi(x) dx \end{aligned}$$

$$\therefore \int \phi(x) dx \neq 0$$

$$\therefore \sum_k h_k = \sqrt{2}$$

**Orthogonality:** For any  $l \in \mathbb{Z}$ ,

$$\sum_k h_k h_{k-2l} = \delta_l$$

$$\begin{aligned} Pf : \quad \phi(x)\phi(x-l) &= \sqrt{2} \sum_k h_k \phi(2x-k)\phi(x-l) & (3.18) \\ &= \sqrt{2} \sum_k h_k \phi(2x-k) \sqrt{2} \sum_m h_m \phi(2(x-l)-m) \end{aligned}$$

By integrating the both sides in (3.18) we obtain

$$\begin{aligned}\delta_l &= 2 \sum_k h_k \left[ \sum_m h_m \frac{1}{2} \int \phi(2x-k) \phi(2x-2l-m) d(2x) \right] \\ &= \sum_k \sum_m h_k h_m \delta_{k, 2l+m} = \sum_k h_k h_{k-2l}\end{aligned}$$

$$\therefore \delta_0 = \int \phi(x) \phi(x) dx = \langle \phi(x), \phi(x) \rangle = \|\phi(x)\|^2 = 1$$

$$\therefore \text{If } l=0, \text{ then } \delta_0 = \sum_k h_k^2 = 1$$

The fact that the system  $\{\phi(\bullet - k), k \in \mathbb{Z}\}$  constitutes an orthonormal basis for  $V_0$  can be expressed in the Fourier domain in terms of either  $\Phi(w)$  or  $m_0(w)$

(a) In terms of  $\Phi(w)$ : 
$$\sum_{l=-\infty}^{\infty} |\Phi(w + 2\pi l)|^2 = 1$$

Pf: Parseval property: 
$$\langle f, g \rangle = \frac{1}{2\pi} \langle \hat{f}, \hat{g} \rangle$$

By the [PAR] property of the Fourier transformation and the  $2\pi$ -periodicity of  $e^{i\omega k}$

$$\begin{aligned}
\delta_k &= \int_R \phi(x) \overline{\phi(x-k)} dx \\
&= \frac{1}{2\pi} \int_R \Phi(w) e^{iwx} \overline{\Phi(w)} e^{-iw(x-k)} dw \\
&= \frac{1}{2\pi} \int_R \Phi(w) \overline{\Phi(w)} e^{iwk} dw \\
&= \frac{1}{2\pi} \int_0^{2\pi} \sum_{l=-\infty}^{\infty} |\Phi(w+2\pi l)|^2 e^{iwk} dw \\
&= \frac{1}{2\pi} \int_0^{2\pi} f(w) e^{iwk} dw
\end{aligned} \tag{3.21}$$

choose  $k = 0$

$$\delta_0 = 1 = \frac{1}{2\pi} \int_0^{2\pi} f(w) dw$$

$$\Rightarrow f(w) = \sum_{l=-\infty}^{\infty} |\Phi(w+2\pi l)|^2 = 1$$

(b) In terms of  $m_0$ :  $|m_0(w)|^2 + |m_0(w + \pi)|^2 = 1$  (3.23)

Pf: Since  $\sum_{l=-\infty}^{\infty} |\Phi(2w + 2l\pi)|^2 = 1$ , then by  $\Phi(w) = m_0(\frac{w}{2}) \Phi(\frac{w}{2})$

$$\sum_{l=-\infty}^{\infty} |m_0(w + l\pi)|^2 |\Phi(w + l\pi)|^2 = 1 \quad (3.24)$$

$$1 = \sum_{k=-\infty}^{\infty} |m_0(w + 2k\pi)|^2 |\Phi(w + 2k\pi)|^2 \\ + \sum_{k=-\infty}^{\infty} |m_0(w + (2k + 1)\pi)|^2 |\Phi(w + (2k + 1)\pi)|^2$$

since the period of  $m_0(w)$  is  $2\pi$

$$\begin{aligned} 1 &= |m_0(w)|^2 \sum_{k=-\infty}^{\infty} |\Phi(w + 2k\pi)|^2 \\ &\quad + |m_0(w + \pi)|^2 \sum_{k=-\infty}^{\infty} |\Phi(w + \pi) + 2k\pi|^2 \\ &= |m_0(w)|^2 + |m_0(w + \pi)|^2 \end{aligned}$$

### Remark 3.3.1

By identity (3.20), any set of independent functions spanning  $V_0$   $\{\phi(x-k), k \in \mathbb{Z}\}$  can be orthogonalized in the Fourier domain. The orthonormal basis is generated by integer-shifts of the function

$$F^{-1} \left[ \frac{\Phi(\omega)}{\sqrt{\sum_{l=-\infty}^{\infty} |\Phi(\omega + 2\pi l)|^2}} \right]$$

This normalization in the Fourier domain is sometimes used in constructing wavelet bases



### **Remark 3.3.2**

The system  $\{\phi(\bullet - k), k \in \mathbb{Z}\}$  is a frame for  $V_0$   
iff

$$A \leq \sum_{l=-\infty}^{\infty} |\Phi(w + 2l\pi)|^2 \leq B$$

Where A and B are frame bounds.

### Remark 3.3.3

Conditions  $|m_0(w)|^2 + |m_0(w + \pi)|^2 = 1$  and  $\sum_{k=-\infty}^{\infty} |\Phi(2w + 2k\pi)|^2 = 1$  are

**not** equivalent. The first is a necessary and the second is a sufficient condition for orthogonality.

*For example,*

$$\phi(x) = \frac{1}{3} \mathbf{I}(0 \leq x \leq 3) \quad \text{has} \quad m_0(w) = \frac{1 + e^{-3iw}}{2}$$

$$\text{It satisfies } m_0(0) = 1 \quad \text{and} \quad \left| \frac{1 + e^{-3iw}}{2} \right|^2 + \left| \frac{1 - e^{-3iw}}{2} \right|^2 = 1$$

$$\text{Since } \Phi(w) = e^{-3iw/2} \frac{\sin 3w/2}{3w/2}$$

$$\sum_{k=-\infty}^{\infty} |\Phi(2w + 2k\pi)|^2 = \frac{1}{3} + \frac{4}{9} \cos w + \frac{2}{9} \cos 2w \neq 1$$

### 3.3.1 Derivation of a Wavelet Function

Whenever a sequence of subspaces satisfies MRA properties, there exists (though *not unique*) an orthonormal basis for  $L_2(\mathbb{R})$

$$\left\{ \psi_{jk}(x) = 2^{j/2} \psi(2^j x - k), j, k \in \mathbb{Z} \right\} \quad (3.25)$$

such that  $\left\{ \psi_{jk}(x), j \text{ fixed}, k \in \mathbb{Z} \right\}$  is an orthonormal basis of the “difference space”  $W_j = V_{j+1} - V_j$

The function  $\psi(x) = \psi_{00}(x)$  is called a **wavelet function** or **informally the mother wavelet**.

Since  $\psi(x) \in V_1$  (because of the containment  $W_0 \subset V_1$ ), it can be represented

$$\text{as } \psi(x) = \sum_{k \in \mathbb{Z}} g_k \sqrt{2} \phi(2x - k) \quad (3.26)$$

for some coefficients  $g_k, k \in \mathbb{Z}$

$$\text{Define } m_1(w) = \frac{1}{\sqrt{2}} \sum_k g_k e^{-ikw} \quad (3.27)$$

By a derivation similar to that in (3.14), we obtain the Fourier counterpart of (3.26),

$$\Psi(w) = m_1\left(\frac{w}{2}\right) \Phi\left(\frac{w}{2}\right) \quad (3.28)$$

Prove  $\sum_{l=-\infty}^{\infty} \Psi(w+2l\pi)\overline{\Psi(w+2l\pi)}=0$  (\*)

Pf:  $\because W_0$  and  $V_0$  are orthogonal,  $\phi(x) \in V_0, \varphi(x) \in V_1$

$$\begin{aligned} 0 &= \int \psi(x)\phi(x-k)dx \\ &= \frac{1}{2\pi} \int \Psi(w)\overline{\Psi(w)}e^{iwk}dw \\ &= \frac{1}{2\pi} \int_0^{2\pi} \sum_{l=-\infty}^{\infty} \Psi(w+2l\pi)\overline{\Psi(w+2l\pi)}e^{iwk}dw \end{aligned}$$

By repeating the Fourier series argument, as in (3.20)

Then we get (\*)

$$(*) \Leftrightarrow m_1(w)\overline{m_0(w)} + m_1(w+\pi)\overline{m_0(w+\pi)} = 0 \quad (3.29)$$

From (3.29), we conclude that there exists a function  $\lambda(w)$  such that

$$(m_1(w), m_1(w+\pi)) = \lambda(w)(\overline{m_0(w+\pi)}, -\overline{m_0(w)}) \quad (3.30)$$

By substituting  $\xi = w + \pi$  and by using the  $2\pi$ -periodicity of  $m_0, m_1$

We conclude that  $\lambda(w) = -\lambda(w + \pi)$  and  $\lambda(w)$  is  $2\pi$ -periodic

**Note:** Any function  $\lambda(w)$  of the form  $e^{\pm iw} S(2w)$ , where  $S$  is an  $L_2([0, 2\pi])$

We choose  $\lambda(w)$  such that

(a)  $\lambda(w)$  is  $2\pi$ -periodic

(b)  $\lambda(w) = -\lambda(w + \pi)$

(c)  $|\lambda(w)|^2 = 1$



Define  $m_1(w) = -e^{-iw} \overline{m_0(w + \pi)}$  (3.32)

Since  $|m_1(w)| = |m_0(w + \pi)|$ , the orthogonality condition (3.23) can be

rewritten as  $|m_0(w)|^2 + |m_1(w)|^2 = 1$

By comparing the definition of  $m_1$  in (3.27) with

$$\begin{aligned}
 m_1(w) &= -e^{-iw} \frac{1}{\sqrt{2}} \sum_k h_k e^{ik(w+\pi)} \\
 &= \frac{1}{\sqrt{2}} \sum_k (-1)^{1-k} h_k e^{-iw(1-k)} \\
 &= \frac{1}{\sqrt{2}} \sum_n (-1)^n h_{1-n} e^{-iwn} \\
 m_1(w) &= \frac{1}{\sqrt{2}} \sum_k g_k e^{-ikw} \tag{3.27}
 \end{aligned}$$

We relate  $g_n$  and  $h_n$  as

$$g_n = (-1)^n h_{1-n} \quad (3.34)$$

In signal processing, the relation (3.34) is known as the **quadrature**

**mirror relation** and the filter  $\tilde{h}$  and  $\tilde{g}$  as **quadrature mirror filters**.

### **Remark 3.3.4**

Choosing  $\lambda(w) = e^{iw}$  leads to the rarely used high-pass filter

$$g_n = (-1)^{n-1} h_{-1-n}$$

It is sometimes convenient to define  $\tilde{g}_n$  as  $\tilde{g}_n = (-1)^n h_{1-n-M}$  where  $M$  is a “shift constant”. Such re-indexing of  $\tilde{g}$  affects only the shift-location of the wavelet function.

*Thanks*